

The rate of facultative sex governs the number of expected mating types in isogamous species

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It is unclear why sexually reproducing isogamous species frequently contain just two self-incompatible mating types. Deterministic theory suggests that since rare novel mating types experience a selective advantage (by virtue of their many potential partners), the number of mating types should consistently grow. However, in nature, species with thousands of mating types are exceedingly rare. Several competing theories for the predominance of species with two mating types exist, yet they lack an explanation for how many are possible and in which species to expect high numbers. Here, we present a theoretical null model that explains the distribution of mating type numbers using just three biological parameters: mutation rate, population size and the rate of sex. If the number of mating types results from a mutation-extinction balance, the rate of sexual reproduction plays a crucial role. If sex is facultative and rare (a very common combination in isogamous species), mating type diversity will remain low. In this rare sex regime, small fitness differences between the mating types lead to more frequent extinctions, further lowering mating type diversity. We also show that the empirical literature supports the role of drift and facultativeness of sex as a determinant of mating type dynamics.

In most sexually reproducing species, gametes do not fuse indiscriminately: syngamy only occurs between gametes of complementary mating types. The evolution of gamete self-incompatibility is puzzling, as self-discriminatory mutants limit their reproductive opportunities. A number of theories for self-incompatibility have been proposed (for a review, see ref. ¹), including modifying the costs of sex by promoting outcrossing, maximizing the rate of attraction between gametes and managing conflict between cytoplasmic organelles.

Suppose that self-incompatibility evolves by one of the selective forces above. The next question that arises is ‘how many mating types might we expect to see?’ A simple answer is very many: any novel type that can fuse with the entire resident population is favoured when rare (negative frequency-dependence) and thus will always invade^{2,3}. Extrapolating this to the extreme predicts a population with as many mating types as individuals. However, such abundances are not observed in the natural world. While examples of species with numerous mating types exist (the fungus *Schizophyllum commune* has over 23,000; ref. ⁴), the vast majority of species feature only 2 (ref. ¹), contradicting our naïve extrapolation above.

A more thorough treatment requires consideration of gamete morphology. The morphological similarity between gametes in isogamous species⁵ contrasts with the clear dimorphism (sperm and eggs) in anisogamous and oogamous species. Explanations for the transition from isogamy (considered the ancestral state^{6–9}) to two sexes (anisogamy) usually consider competition for fertilization driving disruptive selection^{10–12}. If within-sex mating (egg–egg or sperm–sperm syngamy¹³) is detrimental, selection favours linkage between mating types and sexes¹⁴, and it becomes difficult to envisage how a novel mating type could benefit by mating with more than half the population. Conversely, isogamous species should not a priori feature this fundamental restriction; indeed, *S. commune* is isogamous. Yet, the majority of isogamous species have just two mating types¹⁵.

Our aim is to explain the preponderance of species with very few mating types, as well as the existence of species with many more, in isogamous populations. Currently, numerous competing theories exist (for reviews, see refs ^{1,16}), often invoking the same hypothetical mechanisms that generate self-incompatibility. One theory posits that two types maximize the efficiency of pheromone signalling while preventing counterproductive attraction to own pheromones^{13,15,17}. While plausible¹⁸, this hypothesis does not account for exceptions¹⁹ and does not predict precisely when more than two mating types might be expected. Furthermore, the predominance of two types remains somewhat unaccounted for under these models, as the conditions that prevent the emergence of three types appear stringent¹⁸. Alternatively, uniparental inheritance (UPI) of organelles has been suggested to drive the evolution of mating types^{20–22}, with the directionality of inheritance (donor and receiver) limiting the number of mating types. However recent theoretical work accounting for frequency-dependent fitness in UPI has demonstrated that selection for UPI does not obviously lead to the evolution of a low number of mating types²³. A further issue is that the hypothesis does not address the mating type numbers in species without UPI. A similar lack of predictive power applies to the idea that evolutionary constraints on genome architecture might make it difficult for a new type to arise that can mate with others^{24,25}.

Verbally, it has been argued that genetic drift may limit the number of mating types². Here, we model the relevant population genetics. Crucially, we insert an overlooked biological component to previous verbal ideas^{1,2}: we explicitly take into account that the majority of isogamous sexual organisms engage in facultative sex. We show that the consequent alternation of asexual and (rare) sexual cycles has a major effect on mating type dynamics. We derive the expected number of mating types when genetic drift, causing extinction of types, is balanced by mutations, which introduce new mating types to the population.

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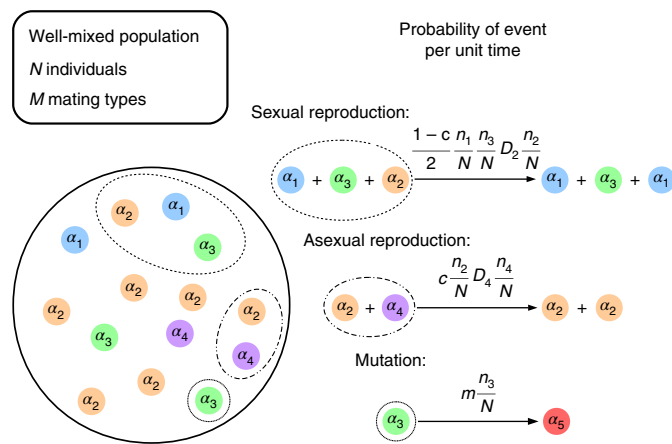


Fig. 1 | Visualization of the model illustrating three types of potential event: sexual reproduction, asexual reproduction and mutation. Each occurs with a probability rate proportional to the frequency of each type involved in the event. In this example, type α_2 is less likely to reproduce sexually than the other types due to its high frequency. However, as it can engage in asexual reproduction, its frequency can still potentially increase due to drift. At a given time, there are M mating types present in the population. All events (a total of $M^2(M-1)$ for sexual reproduction, M^2 for asexual reproduction and M for mutation) lead to one of M^2 different transitions in which one mating type increases and another decreases by 1. Summing over all events that lead to each transition yields equations (1) and (2).

Results

Model. Our model aims to be general by being simple. However, we contextualize it with model isogamous organisms, such as the single-celled green alga *Chlamydomonas reinhardtii* and *Saccharomyces* yeasts. Consider *Chlamydomonas*. In the wild, *C. reinhardtii* mostly exists in a haploid state, replicating asexually through mitosis. Falling nutrition levels instigate the sexual phase of its life cycle²⁶; facultative sex under stress is common across lower eukaryotes²⁷. The haploid cell mitotically divides into four gametes^{28,29}. One of two alleles at a single locus determine the mating type³⁰. Opposite mating types engage in syngamy. Following meiosis, the cells divide, with half inheriting the mating type of each parent.

We take a population genetics approach and use a Moran-type model with a constant population size N (birth–death events are coupled) and overlapping generations³¹. Self-incompatible mating types are determined by an allele at a single locus. To explore the dynamics of mating type numbers, we allow for an infinite number of mating type alleles at this locus. For the full approach, see the Methods; here, we review the salient points.

We denote by α_i mating type i , with number of individuals n_i and frequency $x_i = n_i/N$. The number of mating types present in the population is denoted M . Individuals can experience three classes of event: asexual reproduction, sexual reproduction and mutation. Both types of reproductive event produce a single progeny; for sexual reproduction, the progeny inherits either parent's genotype with probability 1/2 (see Fig. 1). The parameter c controls the rate of sex from $c=1$ (entirely asexual reproduction) to $c=0$ (obligate sex). Our assumption of a constant propensity for sexual reproduction $(1-c)$ means we consider the many mechanisms for the evolution of recombination rates³² beyond the scope of our model. For instance, in the *Chlamydomonas* example, c can be interpreted as the probability that an individual is in a non-stressed state and reproduces asexually.

Mutants arrive independently at a rate m , are novel to the population, self-incompatible, and mate with resident types at the same rate as resident self-incompatible types mate with each other. This

liberal assumption does not account for maladaptation in signaling or syngamy of the mutant with its self-incompatible ancestor. However, since identical fitness across all mating types is an unlikely scenario³³ (see also ref. ³⁴ in the context of sex-ratio evolution), we consider an extension of our main model, in which each type has its own mortality rate, D_i , drawn chosen from a normal distribution with mean 1 and variance σ .

Since each event involves replacing one individual with another, we combine multiple events into a single term, T_{ij} , the probability per unit time of n_i increasing by 1 and n_j decreasing by 1:

$$T_{ij} = \left[\text{Asexual reproduction} \frac{n_i}{cN} + \left(\frac{1-c}{2} \right) \frac{n_i}{N} \left(\frac{\sum_{k \neq i} n_k}{N} \right) \right] \left[\frac{n_j}{D_j N} \right] \text{ if } n_i > 0 \quad (1)$$

$$T_{ij} \propto \left[\frac{n_j}{mN} \right] \text{ if } n_i = 0 \quad (2)$$

Our model includes ideas from Fisherian sex-ratio theory (common types have lowered mating success) while relaxing the assumption, inherent in classic Fisherian theory, that failure to engage in sex implies complete reproductive failure (here, asexuality is still an option). The lowered mating success of common types is taken into account by the term $\sum_{k \neq i} n_k$, the sum over the numbers of all α_k mating types that are not identical to the focal reproducing mating type, α_i ; sex can only occur between different types. Less common mating types will have an increased per-capita probability of participating in sexual events compared with more common types, while there is no difference in the context of asexual reproduction. The strength of selection against common types (negative frequency-dependence) therefore depends on the frequency of sex.

Model behaviour. We first consider an infinite population with rare mutations ($m \ll 1$) and no differential mortality ($\sigma=0$). Denoting by M_0 the initial number of mating types present, the population approaches a state where all types are equally represented ($x_i = 1/M_0$; see Fig. 2a). It resides here until mutation introduces a novel mating type. A new stable state then emerges at $x_i = 1/(M_0 + 1)$. This pattern leads to linear growth in the number of mating types (see equation (10)) and predicts, at very long times, infinitely many mating types each at infinitely low frequency (see Fig. 3a), in agreement with earlier simpler models². Next, we allow differential mating type mortality ($\sigma > 0$). Types are no longer equally represented (see equation (11) and Fig. 2a,b). The departure from even-type frequencies increases with M (the number of mating types), c (the rate of asexual reproduction) and σ (mortality variance). The probability that this polymorphic equilibrium is stable also decreases with these parameters (see Fig. 3b), limiting the number of mating types to a large but finite value.

Turning to finite populations, an infinite number of mating types becomes obviously impossible, even in the absence of mortality differences between the mating types. The number is instead determined by a balance between mutation and extinction. Low mutation rates (that limit the supply of new types) and population sizes (that increase drift) reduce M . These relatively obvious effects co-occur with the more interesting effect of a non-zero propensity for asexual reproduction, c . High c can greatly amplify the effect of drift, which, by speeding up the extinction rate, leads to fewer mating types (see Fig. 3c). Differential mortality rates exacerbate

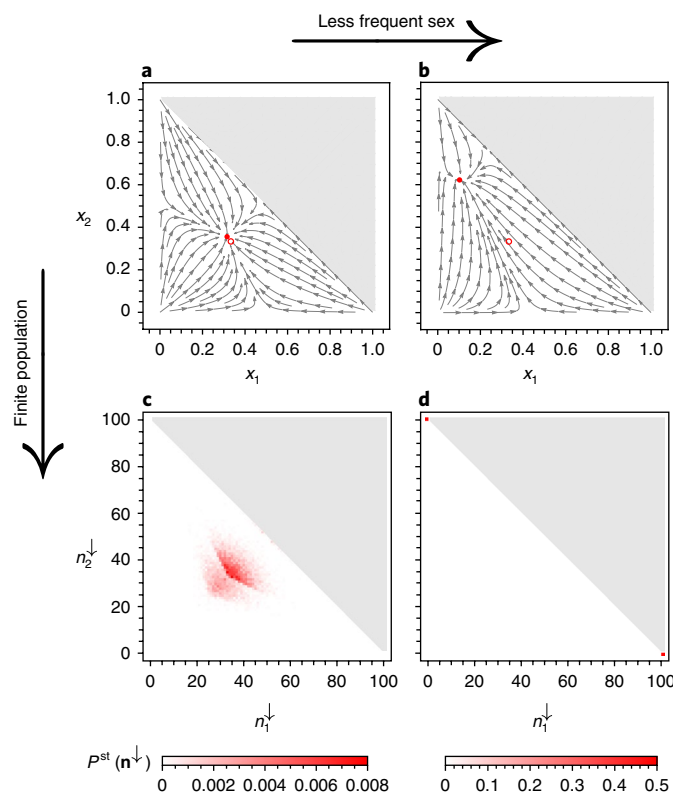


Fig. 2 | Dynamics of mating type frequency with three initial mating types ($M_0 = 3$). **a, b**, Phase plots of the dynamics at the deterministic limit ($N = \infty$; see equation (8)). **c, d**, Results of stochastic Gillespie simulation of equation (2) ($N = 10^3$), with the stationary distribution $P^{\text{st}}(\mathbf{n})$ projected into the two-dimensional plane. In the presence of weak selection ($\sigma = 0.04$, **a–d**), the fixed point of the deterministic dynamics (**a** and **b**, red disk; see equation (11)) is shifted from the neutral prediction $x_i = 1/M_0$ (red open circles). This shift is greater when sex is less frequent ($c = 0$ in plots **a** and **c**; $c = 0.8$ in plots **b** and **d**). In **c** and **d**, drift allows the system to move away from the deterministic fixed point. When sex is rare, both drift and weak selection increase the extinction rate (see **d**, where only a single mating type allele is present at long times). In plots **c** and **d**, $m_g = 5 \times 10^{-3}$.

this process by shifting the equilibrium deterministic fixed point closer to the extinction boundaries (see Fig. 3d).

The above evokes very long-term arguments, and we seek an analytic characterization of the population's behaviour once any transient dynamics have died away. The mathematical analysis is conducted under the assumption that there are no fitness differences between mating types ($\sigma = 0$); the dynamics outside this regime is explored via simulation.

Denote by $\mathcal{P}^{\text{st}}(M)$ the probability of observing M mating types in the population at infinite time and assume no differential mortality. Employing approximations that rely on biologically reasonable assumptions of large N and small per generation mutation rate, $m_g = mN$ (novel mating types arise far less than once each generation), we can obtain bounds on the mode of $\mathcal{P}^{\text{st}}(M)$ under three different scenarios in which all mating types have the same mortality; that is, for obligately sexual ($c = 0$), obligately asexual ($c = 1$) and facultatively sexual ($1 > c > 0$) organisms. For obligate sex, we find

$$\sqrt{\frac{N}{W\left[\frac{1}{4e^2m^2N}\right]}}} > \text{Mode}[\mathcal{P}^{\text{st}}(M)] > \sqrt{\frac{-N}{2(1+\log[2m])}}} \quad (3)$$

where $W[z]$ is the Lambert W function and e is Euler's number³⁵. Thus, for obligately sexual isogamous organisms, we predict many mating types of the order hundreds (see Fig. 4).

For obligately asexual organisms, we find

$$\text{Mode}[\mathcal{P}^{\text{st}}(M)] \approx 1 \quad (4)$$

that is, we expect only a single, non-expressed (in the absence of sex), mating type. If mutants arise much less often than once per generation (our assumption), genetic drift purges these neutral variants from the population faster than they are produced.

If sex is facultative (the most important category for isogamous organisms), the expected number of mating types in the population lies between the upper bound in equation (3) and the lower bound in equation (4). A precise solution can be obtained numerically (see equation (17) and Supplementary Fig. 8), leading to our most important finding (Fig. 4): at low but non-zero rates of sexual reproduction ($1 - c$), large populations can maintain distinct mating types (unlike the asexual system), but the number of mating types can be very low (unlike the sexual system).

Mortality differences ($\sigma > 0$) prevent us from obtaining analytical results of the type described above. Simulating the model, we find that $\sigma > 0$ decreases the expected number of mating types for all values of c (see Supplementary Information). This decrease is larger when c is large (sex is rare) as this leads to more distorted mating type frequencies (see equation (11)) and consequently more frequent extinctions.

Empirical support for the model. We predict low numbers of mating types to associate with a small effective population size (N_e), low mutation rates and rare sex. Before evaluating the relevant evidence, we first list why estimating these parameters is challenging.

Estimating mutation rates producing new mating types is particularly difficult, as we are unaware of a single documented case of a de novo mating type arising. Thus, like Wright³⁶, we simply discuss what is reasonable.

The second parameter, N_e , typically falls below the census population size as various processes (local population structure³⁷, bottlenecks, and sexes and recombination themselves³⁸) can accentuate genetic drift. Although species-specific estimates of N_e vary, we present results based on N_e ranging between 10^6 and 10^7 —a compromise that is both evidence based^{39,40} and avoids interpreting model performance too optimistically (for example, stress-induced sex coinciding with population bottlenecks could lower N_e , yielding a better match between model predictions and observations).

Finally, many otherwise well-studied species are data deficient for their rate of sex. Estimates based on molecular methods⁴¹ are disputed⁴² and data on wild populations^{43,44} are rare. Sex occurs in many organisms previously thought to be asexual⁴⁵, and cryptic sex is common⁴⁶. We focus on well-studied species out of necessity, noting that this may overestimate the frequency of sex: organisms may become model species precisely because sex is straightforward to induce in laboratory settings. When evaluating the propensity for sex in our model ($1 - c$) empirically, a difficulty is that we predict lower c to yield more mating types (higher M); simultaneously, higher M can lead to more mating opportunities and thus amplified signatures of sex. However, disentangling these features is possible. Increasing the number of mating types should lead to modest changes in the amount of sex (matings are possible between 50 and 100% of the population), with larger differences reflecting actual differences in the propensity for sex, ($1 - c$). We proceed with the above caveats in mind.

Consider two closely related yeasts *Saccharomyces cerevisiae* and *S. paradoxus*. Both have two mating types. Molecular studies suggest that sex is exceedingly rare, with estimates of an outcrossing event once per 5,000 asexual divisions⁴⁷ or sex once in every

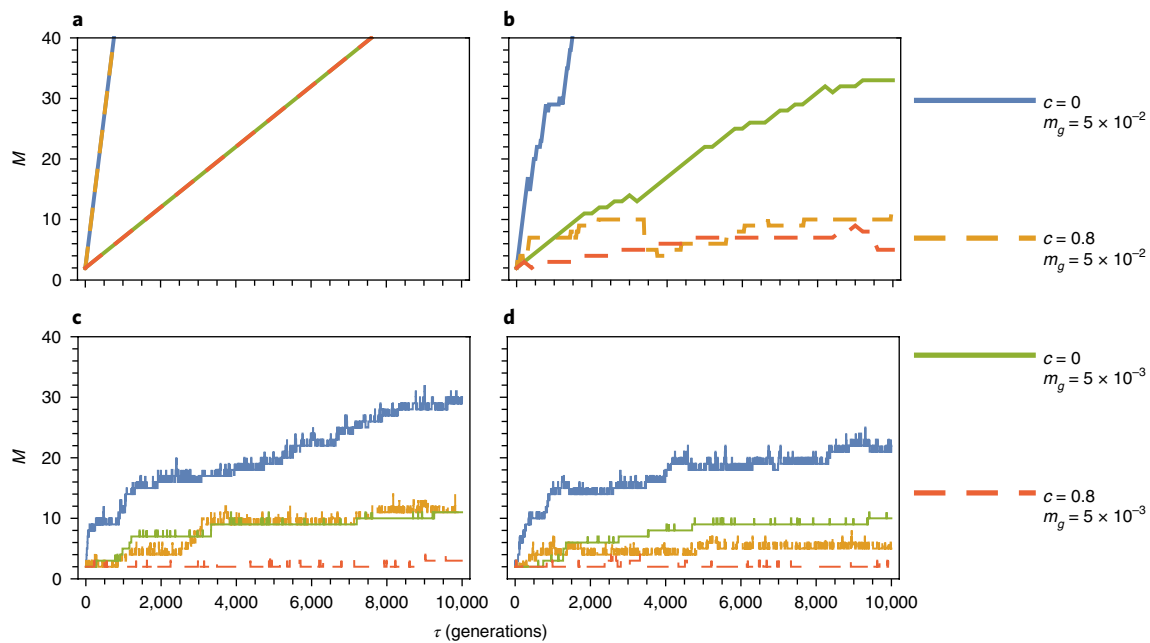


Fig. 3 | Dynamics of the mating type number, M , with $M_0 = 2$ for various parameters under different modelling assumptions. **a–d, Results for dynamics at the deterministic limit ($N = \infty$; **a** and **b**) and stochastic limit ($N = 10^5$; **c** and **d**) for neutral dynamics ($\sigma = 0$; **a** and **c**) and non-neutral dynamics ($\sigma = 0.04$; **b** and **d**). Under the infinite population size limit when all mating types are equally fit (**a**), the rate of asexual-to-sexual reproduction, c , does not affect the dynamics of M (see equation (10)). When weak fitness differences between the alleles are present (**b**), rare sex (large c) decreases the strength of selection for even mating type numbers (see Fig. 2), leading to lower values of M (dashed lines). Finite population sizes (**c** and **d**) further lower mating type numbers through drift-induced extinctions. In this context, lower mutation rates also limit the observed number of mating types.**

1,000 to 3,000 generations⁴¹. For *C. reinhardtii*, molecular studies suggest that just 1,000 outcrossing events may have occurred in isolates sampled over the past 70 years⁴⁸. Assuming $(1 - c) = 1/1,000$ ($c = 0.999$), $10^7 \geq N_e \geq 10^6$ and $10^{-6} \geq m_g \geq 10^{-8}$ (a new mating type once every 1 to 100 million generations), the model predicts $13 \geq M \geq 3$ in the absence of viability selection treating the mating types unequally (see Fig. 4). Although this exceeds the two types observed in *Saccharomyces* and *Chlamydomonas*, our model shows that small mortality differences (of the order 0.1%) can reduce the number of mating types to 2 (see Supplementary Section 5; also see below for other potential factors, such as species-specific genetic architecture).

Ciliates appear to reproduce sexually more often, but still infrequently since sex is limited by an immature period of 40–100 asexual divisions^{19,43}. Studies on wild populations^{43,44} estimate sex to occur once in every few hundred generations. We predict that more frequent sex leads to more mating types. Assuming $10^7 \geq N_e \geq 10^6$ and $10^{-6} \geq m_g \geq 10^{-8}$, but now setting $(1 - c) = 1/200$, the model predicts $28 \geq M \geq 6$ (see Supplementary Information). Known numbers of mating types are 5–13 for *Euplotes*, 2 for *Aspidisca* and 3–9 for *Tetrahymena*^{19,49}.

Most species of mushroom-forming fungi *Agaricomycotina* (a subdivision of the *Basidiomycota*) are obligately sexual⁵⁰. *S. commune*, ($M > 23,000$) belongs to this family, and molecular evidence suggests that it is one of the most sexual species in the fungal kingdom⁴², with high mutation rates³⁹. The fungal kingdom, in general, offers evidence of N_e covarying positively with numbers of mating types¹⁶.

Focusing on *S. commune* and assuming obligate sex, a large effective population size ($N_e = 10^7$) and large mutation rate ($m_g = 10^{-6}$), we predict $520 \geq M \geq 420$ (see Fig. 4), which is well below the 23,000 mating types known to exist. *S. commune* offers good biological reasons for our model underestimate. We assumed a single mating type locus. *S. commune* has tetrapolar mating type determination, with each type defined by two loci, each with two weakly recombining

regions⁴. A mating type is not extinct when its genotype frequency reaches zero, but when one of its mating type alleles is lost. Extinctions become less likely in a system where the mating type allele is carried by many more individuals than the genotype.

Generally, multiple loci are expected to stabilize multiple mating type systems better than multiple alleles at a single locus^{24,25,51}. Multiple loci indeed frequently determine isogamous species' mating types, with the gain or loss of loci causing mating type number transitions (for example, *Paramecium bursaria*¹⁹, *Ustilago hordei*, *Malassezia globosa* and *Cryptococcus neoformans*⁴⁵). Single-locus determination of more than two mating types is only common when mating types are determined at the diploid stage. The disassociation between mating type alleles and mating types then resembles that of multi-locus systems.

More drastic departures from our theory are possible due to genetic architecture. While the highly sexual *S. commune* has thousands of mating types, other highly sexual fungal species have abandoned mating types altogether. Homothallism in *Ascomycetes*⁵² makes them lack bifactorial mating type determination, potentially limiting the scope for novel mating types²⁴. Aligning homothallism with our model is possible if high rates of sex (observed in these taxa, suggesting an ecology that favours high recombination) would permit large numbers of mating types, but this route is limited by mutation. Mutations for self-compatibility then offer an alternative route to maintain frequent sex.

Turning to frequency-independent success differences between mating types, the model predicts that the signals of such differences should manifest more strongly when sex is rare. Consequently, these signals are best sought during periods of asexual reproduction. In *Chlamydomonas*, bouts of asexuality frequently sweep single mating types to fixation⁵³, as a result of hitch-hiking on beneficial mutations and asexuality maintaining the linkage. Similar dynamics occur in fungi^{54,55}, where among pathogenic species^{56–58} there is also evidence for fitness differences between mating type alleles⁵⁹.

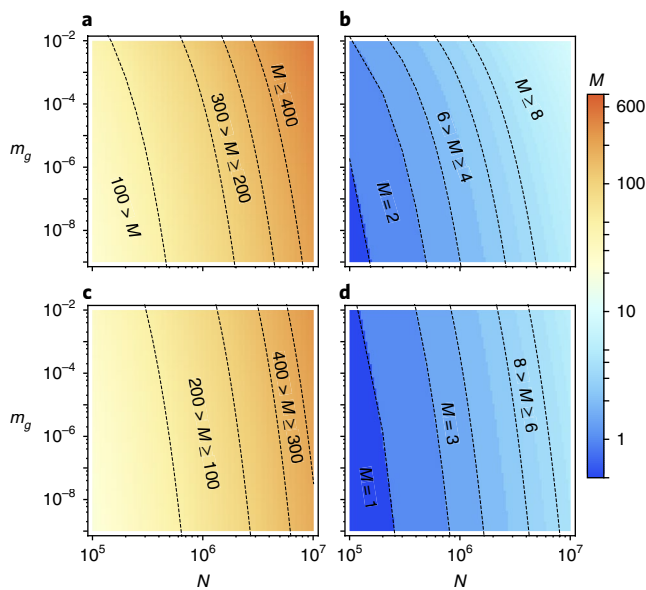


Fig. 4 | Bounds on the expected number of mating types predicted by the model as a function of the population size, N , and the per-generation mutation rate, m_g , when there are no selective differences between the mating types ($\sigma = 0$). a–d, Upper (a and b) and lower bounds (c and d) are shown for an obligately sexual population ($c = 0$; a and c) and a facultatively sexual population in which sex is rare ($c = 0.999$; b and d). Bounds were obtained by evaluating equation (3) in a and c, and by numerically solving equation (17) in b and d.

Our model could be criticized where it predicts two mating types for lacking a mechanism, such as obligate sex in Fisherian sex-ratio theory, that would prevent strong fluctuations around a 1:1 mating type ratio. Here, the parameter regimen places the system precariously near the loss of all but one of the types. We consider these dynamical features real rather than a flaw. A common type is bound to have greatly diminished reproductive success under Fisherian dynamics with sex assumed obligate. However, facultativeness of sex means it can still reproduce (for other routes to deviations from 1:1, see refs ^{34,60}). Balancing selection is not always sufficient to maintain equal mating type ratios in *Coccidioides*⁶¹ ($M = 2$) or *Dictyostelium discoideum*⁶² ($M = 3$). In ciliates, at least two species of *Tetrahymena* have lost mating types¹⁹ (*T. ellioti* and *T. pyriformis*). In the fungal kingdom, *Batrachochytrium dendrobatidis* only has a single mating type allele⁴⁵. While contributions of drift and selection are hard to disentangle for each case, our model suggests that they may act synergistically to reduce the number of mating types.

Discussion

Why do isogamous species have few (and species-specific numbers of) mating types, when the naïve prediction is that rare types should always invade? Our model contributes towards understanding this discrepancy through two important interacting ingredients: finiteness of populations (genetic drift) and facultative sex. These change the prediction from unbounded increases in the number of mating types (the prediction in infinite populations) to a species-specific number of mating types. This number is reduced if mutations yielding new mating types are rare, populations are small and sexual reproduction is rare (compared with the number of asexual cycles).

Our model derives precise expectations for the effects of drift (discussed by refs ^{1,2,62}) across scenarios that differ in their rates of sex and effective population sizes, when all mating types are equally fit and when they are not. When equally fit, the model is a ‘null model’ with no ecological differences beyond frequency-dependent selection favouring rare types. Assessing the role of drift in populations

that can undergo asexual as well as sexual cycles^{63,64} necessitates a stochastic (rather than deterministic^{2,3}) modelling approach. We derive the null expectations without having to rely on simulations⁶³, while avoiding the mathematical inconsistencies of earlier studies of plant self-incompatibility systems⁶⁵.

Our generalized model includes frequency-independent fitness differences between mating types. Here, simulations show that small fitness differences are sufficient to further reduce the number of mating types when sex is rare. Note that the mechanisms incorporated in our model and others (non-mass-action mating kinetics², pheromone signalling^{15,17} and organelle inheritance²⁰) are not mutually exclusive. We expect drift to be important because isogamous sex is typically facultative, and this potentially enhances any mechanism evoked to constrain the success of rare types. The stronger fluctuations in mating type ratio caused by facultative (rather than obligate) sex also help explain the evolution of solutions to mate-finding difficulties; that is, mating type switching or homothallism^{52,66}. Incorporating drift under facultative sex in models devoted to understanding the significance of such mechanisms appears fruitful: low rates of sex may, for example, decrease the costs of biparental inheritance (vegetative segregation generates homoplasmy) and amplify mate-finding problems⁶⁷.

Our model allows us to paint the following picture. Ecological conditions select for high or low rates of sex in a given facultatively sexual species^{68,69}. All else being equal, higher rates reduce drift, permitting more mating types to coexist at an evolutionary equilibrium. While available data fit our model qualitatively, other factors will play a role in the diversity of sexual strategies across taxa. We have discussed homothallism, but another route to a fundamentally different arrangement is male–female dimorphism (anisogamy), in which the number of sexes will be two based on different processes from the ones we envisage. Anisogamy models routinely produce only two size classes⁹ (sperm and egg), and as it is advantageous to prevent sperm attempting fusing with other sperm (as neither gamete would provide sufficient cytoplasm for future development), it is logical to suspect that mating types become associated with size-based classifications of gametes (for molecular evidence, see ref. ⁷).

Methods

For notational convenience, we initially set $M_{\max}$ as the maximum number of possible mating types. The vector $\mathbf{n}$, which describes the number of individuals of each mating type, is then of length $M_{\max}$. We denote by $T(\mathbf{n}'|\mathbf{n})$ the probability per unit time of transitioning from a state $\mathbf{n}$ to state $\mathbf{n}'$. In general, the probability $P(\mathbf{n}, t)$ of being in a state $\mathbf{n}$ at time t is given by

$$\frac{dP(\mathbf{n}, t)}{dt} = \sum_{\mathbf{n}'} [T(\mathbf{n}|\mathbf{n}')P(\mathbf{n}', t) - T(\mathbf{n}'|\mathbf{n})P(\mathbf{n}, t)] \quad (5)$$

To align this mathematical formulation with the model described in the main text, we must define the term $T(\mathbf{n}|\mathbf{n}')$, which in the Moran model consists of the birth of one individual and death of another (see equations (1) and (2));

$$T_{ij} \equiv T[(\dots, n_i + 1, \dots, n_j - 1, \dots)](\dots, n_p, \dots, n_p, \dots) \\ = c \frac{n_i D_j}{N} \frac{n_j}{N} + \left(\frac{1-c}{2} \right) \frac{n_i}{N} \left(\frac{N-n_i}{N} \right) D_j \frac{n_j}{N} \quad \text{if } n_i > 0 \quad (6)$$

$$= \left(\frac{m}{M_{\max} - M} \right) \frac{n_j}{N} \quad \text{if } n_i = 0 \quad (7)$$

Here, M is the number of mating types present in the population (that is, the number of non-zero entries in $\mathbf{n}$) and the term $1/(M_{\max} - M)$ is a normalization factor that accounts for the fact that a new mutation may be assigned to any of the unoccupied mating type labels. Note that we have used the property that $\sum_{i=1}^{M_{\max}} n_i = N$ to simplify equation (1) to (6). Also note that in the limit $c \rightarrow 1$, the model simplifies to the neutral infinite allele Moran model with mutation⁷⁰. A similar modelling approach has been used to investigate the number of self-incompatibility alleles in plants^{36,63,71}, but focused on diploid systems and without accounting for asexual reproduction. Meanwhile our generic mating kinetics (mass action) is the same as that used in the first of four models explored in ref. ².

The time for N reproduction events is approximately N in units of t (see Supplementary Section 2). Therefore, we introduce τ as the generation time,

$\tau = t/N$. In a similar fashion, we introduce $m_g = mN$ as the per-generation mutation rate. Our model does not account for the possibility that a gamete chosen for sexual reproduction fails to find a mate, unlike the models described in ref. ² (see the sections ‘Mating Kinetics 2’ and ‘Mating Kinetics 3’ within this reference). Finally, note that equation (6) implicitly assumes that sexual reproduction events are not temporally correlated. For an alternative approach, see ref. ⁶⁶.

To obtain the deterministic (infinite population size) limit to the dynamics, we can apply a diffusion approximation⁷⁰. Assuming large N , we transform into the approximately continuous variables $x_i = n_i/N$, Taylor expand equation (5) in N^{-1} and take the limit $N \rightarrow \infty$ ⁷². Recall that the values for D_i are chosen from a normal distribution with mean 1 and s.d. σ . We can therefore rewrite these terms as $D_i = 1 + \sigma d_i$, where d_i is normally distributed with mean 0 and s.d. 1. Assuming that σ is small and the mutation rate is much smaller, we obtain the description (see Supplementary Section 2)

$$\begin{aligned} \frac{dx_i}{d\tau} = & \frac{(1-c)}{2} \left[\sum_{j=1}^{M_{\max}} x_j x_i (x_j - x_i) \right] \\ & + \sigma \sum_{j=1}^{M_{\max}} x_j x_i \left[\left(\frac{1-c}{2} \right) (d_j x_j - d_i x_i) \right. \\ & \left. + \left(\frac{1+c}{2} \right) (d_j - d_i) \right] \end{aligned} \quad (8)$$

This dynamics, illustrated in Fig. 3, recapitulates that of the section ‘Mating Kinetics 1’ in ref. ² when $c=0$ and $\sigma=0$.

When $\sigma=0$, given an initial number M_0 of mating types present in the population, a fixed point exists at $x_i = 1/M_0$ for the present types and $x_i=0$ otherwise. This situation represents even mating type ratios. Considering just the present mating types, the fixed point has eigenvalues

$$\lambda_i = -\frac{1-c}{2M_0} \quad (9)$$

and is therefore stable as long as there is at least some sexual reproduction; that is, $c < 1$ (recall that if $c=1$ the model becomes the neutral infinite allele Moran model, in which the dynamics is governed entirely by drift). If a novel mating type is introduced into the population, this fixed point becomes unstable, and a new fixed point arises at $x_i = 1/(M_0 + 1)$. The stability of the fixed point decreases with increasing c and M_0 . However, since it is always stable for $c < 1$, the number of mating types is expected to grow linearly with time;

$$M(\tau) = M_0 + m_g \tau \quad (10)$$

If $\sigma > 0$, a fixed point at even sex ratios is no longer possible. The fixed point is dependent on the stochastically chosen values of d_i , and is given by

$$x_i = \frac{1}{M_0} + \sigma \frac{1-c-M_0-cM_0}{(1-c)M_0^2} \left(M_0 d_i - \sum_{j=1}^{M_{\max}} d_j \right) \quad (11)$$

for each type present in the population. Note that when c , M_0 and σ are small (frequent sex, a low number of mating types and small variance in mating type fitness), the deviation from even sex ratios is small. However, as c , M_0 and σ increase, deviations from even sex ratios become more pronounced. This additionally leads to an increased probability that the interior fixed point containing all M_0 mating types becomes unstable (or disappears entirely); however, we do not analytically quantify this effect here.

Next, we aimed to analytically characterize the probabilistic dynamics. For this, we wished to obtain the stationary probability distribution $P^{\text{st}}(\mathbf{n})$, which is the solution to equation (5) in the limit $t \rightarrow \infty$. In the Supplementary Information, we show that this equation can be solved, but only if there are no frequency-independent fitness differences between the mating type alleles ($D=1$ or equivalently $\sigma=0$). Under these conditions, the probability transition rates, equations (6) and (7), can be decomposed into products of birth and death rates that depend only on the number of individuals belonging to the mating types that increase and decrease;

$$T[(\dots, n_i + 1, \dots, n_j - 1, \dots) | (\dots, n_p, \dots, n_q, \dots)] = b(n_i) d(n_j) \quad (12)$$

where

$$\begin{aligned} b(k) = & \underbrace{c \binom{k}{N}}_{\text{Asexual reproduction}} + \underbrace{\left(\frac{1-c}{2} \right) \binom{k}{N} \left(\frac{N-k}{N} \right)}_{\text{Sexual reproduction}} \text{ if } k \geq 1, d(k) \\ = & \underbrace{\left(\frac{k}{N} \right)}_{\text{Death/mutation}} \text{ for all } k, \\ b(0) = & \underbrace{\left(\frac{m}{M_{\max} - M} \right)}_{\text{Mutation}} \end{aligned} \quad (13)$$

Given the functional form of equation (12), the stationary distribution can be given by (see Supplementary Information)

$$P^{\text{st}}(\mathbf{n}) \propto \prod_{i=1}^{M-1} \prod_{k=0}^{n_i^{\dagger}-1} \frac{b(k) d \left(N - k - \sum_{j=1}^{i-1} n_j^{\dagger} \right)}{b \left(N - k - \sum_{j=1}^{i-1} n_j^{\dagger} - 1 \right) d(k+1)} \quad (14)$$

where $\mathbf{n}^{\dagger}$ is the vector $\mathbf{n}$ with its elements rearranged in descending order and M is the number of non-zero elements of $\mathbf{n}$ (see Supplementary Section 1). In a distinct model set-up, a similar result has been obtained⁷³. As this result also relies on transitions having forms of the type given in equation (12), the expressions may be relatable. Equation (14) perfectly captures the results of simulations (see Supplementary Figs. 2–4). Classic investigations into the number of self-incompatibility alleles in plants also rely on calculating the stationary distribution of alleles. However, they do so having applied a diffusion approximation and consider the interactions between a single focal self-incompatibility allele and a population fixed at some prescribed frequency distribution—an approach criticized by Moran based on its fundamental mathematical inconsistencies (see ref. ⁶⁵ for discussion). This approach was necessary as the models of self-incompatibility alleles in plants feature diploid sex determination, and so transitions do not follow the functional form given in equation (12).

We are interested in the stationary distribution of the number of mating types present in the population $P^{\text{st}}(M)$, which is related to $P^{\text{st}}(\mathbf{n})$ by

$$P^{\text{st}}(M) = \sum_{\mathbf{n} \in S^{\text{st}}(M)} P^{\text{st}}(\mathbf{n}) \quad (15)$$

where $S^{\text{st}}(M)$ is the set of all vectors $\mathbf{n}$ that contain M non-zero elements. Since this expression is unintuitive, we proceed to characterize the mode of $P^{\text{st}}(M)$. The calculation is described in full in Supplementary Section 4. To begin, we note that if N is large, and $m_g = mN$ small, $P^{\text{st}}(\mathbf{n})$ will consist of a series of peaks, each located at states where the frequency of the mating types is approximately given by the deterministic fixed points. The mode of $P^{\text{st}}(\mathbf{n})$ can then be obtained by considering its values at states $\mathbf{n}$ in the proximity of successive fixed points; $n_i \approx N/M$ for M mating types and zero otherwise. Obtaining the mode of $P^{\text{st}}(M)$ is more complicated; it depends on values of $P^{\text{st}}(\mathbf{n})$ far from the fixed points (see equation (15)). However, we can consider limiting behaviour of $P^{\text{st}}(\mathbf{n})$ to calculate the upper and lower bounds of the mode of $P^{\text{st}}(M)$, which simulations tell us typically lies close to the mean (see Supplementary Information). For a lower bound, we assume that $P^{\text{st}}(\mathbf{n})$ is constructed from a series of delta peaks at the deterministic fixed points. For an upper bound, we assume that $P^{\text{st}}(\mathbf{n})$ is completely flat in the region around the deterministic fixed points. The full calculation is detailed in Supplementary Section 4. For a general facultatively sexual system, the upper and lower bounds of the mode of $P^{\text{st}}(M)$ can be given by (see Supplementary Section 4.2.3)

$$M_{\text{UB}}^* > \text{Mode}[P^{\text{st}}(M)] > M_{\text{LB}}^* \quad (16)$$

where M_{UB}^* and M_{LB}^* can be obtained as solutions to the equations

$$R(M_{\text{LB}}^*) = 1, \text{ and } \left(\frac{N}{M_{\text{UB}}^* - 1} - 1 \right) R(M_{\text{UB}}^*) = 1 \quad (17)$$

with

$$\begin{aligned} R(M) = & 2m(1+c)^{-\frac{1}{2} + \frac{1+c}{1-c}N} M^{-1 + \frac{3}{2}M + \frac{1+c}{1-c}MN - N} \times \\ & (M-1)^{-\frac{3}{2}(M-1) + \frac{2-cM-M}{1-c}N} \times \\ & (cM + M - 2)^{\frac{1}{2}(M-1) - \frac{2-cM-M}{1-c}N} \times \\ & (cM + c + M - 1)^{-\frac{M}{2} - \frac{1+c}{1-c}MN + N} \end{aligned} \quad (18)$$

A numerical solution to this equation for a given set of parameters can be obtained quickly using a standard root-finding algorithm. Comparing analytical results with those from stochastic Gillespie simulation⁷⁴ with $\sigma=0$, we find excellent agreement (see Supplementary Figs. 6–8).

Reporting Summary. Further information on experimental design is available in the Nature Research Reporting Summary linked to this article.

Code availability. C++ Code for the stochastic simulations conducted in this paper is available at <https://github.com/gwaconstable/FiniteNMatingTypes>.

Data availability. Data generated during the study are available at <https://github.com/gwaconstable/FiniteNMatingTypes>.

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Author contributions

G.W.A.C. designed the project and conducted the mathematical analysis. G.W.A.C. and H.K. developed the model and wrote the paper.

Competing interests

The authors declare no competing interests.

Additional information

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Data collection

Data was generated using custom simulations written in c++. This code is freely available at <https://github.com/gwaconstable/FiniteNMatingTypes>. Mathematical calculations are conducted by hand but we have also written a Mathematica file through which our calculations can be verified. This file is also available at <https://github.com/gwaconstable/FiniteNMatingTypes>.

Data analysis

Data was analyzed and visualized in Mathematica 11.1.1.0. This file is available at <https://github.com/gwaconstable/FiniteNMatingTypes>

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Study description	This theoretical study of mating type evolution was conducted using mathematical analyses and computer simulation.
Research sample	Data was generated using stochastic Gillespie simulation.
Sampling strategy	Where statistics are used, large samples of simulated data (N>10000) were generated. Where single simulations have been used for illustrative purposes, these were chosen as the first simulation that was run. These can be verified using the data generation files at https://github.com/gwaconstable/FiniteNMatingTypes
Data collection	This has been recorded directly from simulation.
Timing and spatial scale	N/A. Data is simulated.
Data exclusions	No data has been excluded.
Reproducibility	We have provided freely code that enables readers to generate all plots in the study from scratch, including simulated data generation (see https://github.com/gwaconstable/FiniteNMatingTypes).
Randomization	N/A. Data is simulated.
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Reporting for specific materials, systems and methods

Materials & experimental systems

n/a	Involved in the study
☒	☐ Unique biological materials
☒	☐ Antibodies
☒	☐ Eukaryotic cell lines
☒	☐ Palaeontology
☒	☐ Animals and other organisms
☒	☐ Human research participants

Methods

n/a	Involved in the study
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☒	☐ Flow cytometry
☒	☐ MRI-based neuroimaging